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# EXACT SOLUTIONS TO A FAMILY OF FITZHUGH-NAGUMO-TYPE EQUATIONS

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# **Abstract**

A family of partial differential equations of type Fitzhugh-Nagumo is studied in this paper, generalizing a result presented in [7]. After an introduction to the problem, we find the exact solutions of the equation under study. The first integral method is employed.

### 1. Introduction

Partial differential equations of type FitzHugh–Nagumo bear the name from FitzHugh and Nagumo who first studied equations of this type in [5], [9] as a simplified nerve conduction model. Such nonlinear reaction—diffusion equations attracted great interest lately because of its potential application in Medicine. For example, it was used as a mathematical model of cardiac excitation [1]. It is also met in circuit theory, biology and genetics [2]. Another mathematical model related to the heart problems has been proposed in [10]. It is a FitzHugh–Nagumo type system. Recently, the authors in [6] investigate the Tzitzeica—Dodd—Bullough partial differential equation of the form  $u_{xt} = e^{-u} + e^{-2u}$ , using the Exp-method, pointing out new periodic solutions.

The partial differential equation considered in this paper is given by:

$$u_t - u_{xx} = Mu^3 + Nu^2 + Pu + Q (1)$$

with  $M, N, P, Q \in \mathbf{R}$ .

When M=-1, N=Q=0, P=1, the equation (1) becomes the real Newell—Whitehead equation and for  $M=-1, N=\alpha+1, P=-\alpha, Q=0$ ,  $\alpha \in \mathbb{R}$ , it falls on the classical FitzHugh–Nagumo equation [7].

Information about some exact solutions of (1) can be found in [11]. The aim of this paper is to obtain another exact solutions, namely travelling-wave solutions, for (1) using the first integral method. This method was used by the authors in [3], [4] for searching solutions of other nonlinear partial differential equations.

#### 2. Exact Solutions

In order to proceed to the investigation of the equation (1), we look for solutions of the form  $u(x,t)=u(\eta)$ , with  $\eta=x+\beta t,\ \beta\in\mathbf{R}$ . Replacing  $u(x,t)=u\left(x+\beta t\right)=U(\eta)$  in (1) one gets

$$\beta U' - U'' = MU^3 + NU^2 + PU + Q. \tag{2}$$

Denoting U' = V, from (2) a system of differential equations of first order is obtained:

$$U' = V$$

$$V' = \beta V - (MU^3 + NU^2 + PU + Q).$$
(3)

In order to solve (3) the first integral method will be used. For a nontrivial solution  $(U(\eta), V(\eta))$  of (3) we search a first integral of (3) in the form  $f(U, V) = a_0(U) - V$ . Then

$$\frac{\partial f}{\partial \eta} = \frac{\partial f}{\partial U} \frac{\partial U}{\partial \eta} + \frac{\partial f}{\partial V} \frac{\partial V}{\partial \eta} = 0 \tag{4}$$

leads to

$$a_0'(U) \cdot V - \beta V + MU^3 + NU^2 + PU + Q = 0.$$
 (5)

The condition f(U, V) = K leads to  $V = a_0(U) - K$ . In the following computation the constant K does not play a special role, so we will consider K = 0, i.e.  $V = a_0(U)$ .

From (5) the following relation is obtained

$$a_0(U)(a_0'(U) - \beta) + MU^3 + NU^2 + PU + Q = 0.$$
 (6)

By identifying the coefficients of U in (6) it results that  $a_0$  (U) must be a polynomial of second degree,  $a_0$  (U) =  $AU^2 + BU + C$  whose coefficients must satisfy

$$\begin{cases} 2A^{2} + M = 0\\ 3AB - A\beta + N = 0\\ 2AC + B^{2} - B\beta + P = 0\\ BC - C\beta + Q = 0. \end{cases}$$
(7)

Let us denote  $M = -2S^2, S > 0$ .

The unknowns of (7) are A, B, C and  $\beta$ .

By replacing  $A=S,\ S>0$  (similar computation can be made for A=-S), respectively B and C as function of  $\beta$  in the last relation (7) one can determine  $\beta$  from the equation

$$4S^{3}\beta^{3} - 3\beta S\left(N^{2} + 6PS^{2}\right) - 54QS^{4} - 9NPS^{2} - N^{3} = 0.$$
 (8)

The equation (8) has at least a real solution.

For a solution  $\beta_1 \in \mathbf{R}$  of (8) the coefficients of  $a_0(U)$  are given by

$$A = S,$$

$$B = \frac{1}{3S} (-N + S\beta),$$

$$C = -\frac{1}{18S^3} (NS\beta + N^2 + 9PS^2 - 2S^2\beta^2).$$
(9)

In order to determine the travelling wave solution of (1) the relation f(U,V)=0 will be used. It is an ordinary differential equation of first degree:

$$U' = AU^2 + BU + C. (10)$$

We consider first the case

$$\Delta := B^2 - 4 \cdot A \cdot C = \frac{-2S^2\beta_1^2 - 12S^2P + 2N^2}{9A^2} > 0,$$

i.e.  $\beta_1^2 < \frac{N^2 - 6S^2P}{S^2}$  (do not forget that S>0). In this situation the equation  $AU^2 + BU + C = 0$  has two real solutions  $\alpha_1$  and  $\alpha_2$ . The equation (10) becomes

$$U' = A(U - \alpha_1)(U - \alpha_2).$$

Its general solution is

$$U(\eta) = \frac{\alpha_2 - K \cdot \alpha_1 \cdot e^{A(\alpha_2 - \alpha_1)\eta}}{1 - K \cdot e^{A(\alpha_2 - \alpha_1)\eta}}.$$
 (11)

In order to define the solution for every  $x \in \mathbf{R}$  and for every  $t \geq 0$  the condition K < 0 must be imposed. By considering  $K = -\lambda^2$  we obtain the solution

$$u(x,t) = \frac{\alpha_2 + \lambda^2 \cdot \alpha_1 \cdot e^{A(\alpha_2 - \alpha_1)(x + \beta_1 t)}}{1 + \lambda^2 \cdot e^{A(\alpha_2 - \alpha_1)(x + \beta_1 t)}}.$$
 (12)

In the case  $\Delta=0$ , the equation  $AU^2+BU+C=0$  has a double real solution  $\alpha=-\frac{B}{2A}$  and the equation (10) has the general solution

$$U(\eta) = \frac{-1}{A\eta + K} - \frac{B}{2A}.\tag{13}$$

There are not values of K for which the solution is defined for every  $x \in \mathbf{R}$  and for every  $t \geq 0$ , hence in this situation some spatial constraint are imposed, namely  $x < -K - \beta t$  or  $x > -K - \beta t$ .

imposed, namely  $x < -K - \beta t$  or  $x > -K - \beta t$ . For  $\Delta < 0$ , i.e.  $\beta_1^2 > \frac{N^2 - 6S^2P}{S^2}$  the general solution of (10) is

$$U(\eta) = \sqrt{p} \tan \sqrt{p} (A\eta + K) - \frac{B}{2A}, \tag{14}$$

with 
$$p = -\frac{B^2 - 4AC}{4A^2} > 0$$
.

The solution

$$u\left(x,t\right) = \sqrt{p}\tan\sqrt{p}\left(A\left(x+\beta t\right) + K\right) - \frac{B}{2A}$$

can not express the action of a dynamical system because it can not be defined for every  $x \in \mathbf{R}$  and  $t \ge 0$ .

Remark 1. Another method for searching exact solutions for equations of the form  $u_t - u_{xx} = -(u - u_1) (u - u_2) (u - u_3), u_1 < u_2 < u_3$  and  $u_1, u_2, u_3 \in \mathbf{R}$  is presented in [8]. Our approach in this paper allows to find exact solutions for such equations even though  $u_1, u_2, u_3$  are complex numbers (Example 2), which is not presented in [8].

Example 1. Let us consider a specific case

$$u_t - u_{xx} = Nu^2 - u^3,$$

corresponding to M=-1,  $\left(S=1/\sqrt{2}\right),$  P=0, Q=0 and arbitrary  $N\neq 0$  in (1) .

The equation (8) becomes

$$\beta^3 - \frac{3}{2} N^2 \beta - \frac{\sqrt{2}}{2} N^3 = 0.$$

Its solutions are  $\beta_1 = \beta_2 = -\frac{N\sqrt{2}}{2}$  and  $\beta_3 = N\sqrt{2}$ .

In the case  $\beta_1 = -\frac{N\sqrt{2}}{2}$  the values the coefficients involved in the computation are  $A = \frac{\sqrt{2}}{2}$ ,  $B = -\frac{N\sqrt{2}}{2}$ , C = 0 and  $\Delta = B^2 > 0$ . In this case  $\alpha_1 = 0$  and  $\alpha_2 = N$  and (12) gives the analytical form of an exact solution:

$$u\left(x,t\right) = \frac{N}{1 + \lambda^{2} \cdot e^{\frac{\sqrt{2}}{2}N\left(x - \frac{N\sqrt{2}}{2}t\right)}}.$$

For  $\beta = N\sqrt{2}$  we obtain  $A = \frac{\sqrt{2}}{2}$ , B = 0 and C = 0. In this case  $\Delta = 0$  and an other exact solution is given by (13):

$$u(x,t) = \frac{-2}{\sqrt{2}x + 2Nt + 2K}. (15)$$

For a given K the solution is defined on

$$D_1 = \left\{ (x,t) \quad \mid \quad t \geq 0 \right. , \left. x < N\sqrt{2}t - K 
ight\}$$

or on

$$D_2 = \left\{ (x,t) \quad | \quad t \ge 0 \ , \ x > N\sqrt{2}t - K \right\}.$$

**Example 2.** Let us consider now another case corresponding to  $M=-2,\ (S=1),\ N=P=0,\ Q=-2.$  Eq (1) becomes

$$u_t - u_{xx} = -2u^3 - 2.$$

In this case (8) leads to  $\beta^3 + 27 = 0$ , having the real root  $\beta = -3$ . We have also A = 1, B = -1, C = 1. From (14), the solution is given by:

$$u(x,t) = \sqrt{p} \tan \sqrt{p} (A(x+\beta t) + K) - \frac{B}{2A}$$

i.e.

$$u(x,t) = \frac{\sqrt{3}}{2} \tan \frac{\sqrt{3}}{2} (K - 3t + x) + \frac{1}{2}.$$

## 3. Conclusions

In this work we investigated a family of partial differential equations of type Fitzhugh—Nagumo, with a general cubic function on the right hand side of the equation, generalizing a result presented in [7]. Exact solutions of this family of equations are pointed out. As applications are presented two examples.

# 4. Acknowledgements

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## References

- [1] R.R. Aliev, A.V. Panilov: A simple two-variable model of cardiac excitation, Chaos Solitons and Fractals, Vol. 7, (1996), 293-301.
- [2] D.G. Aronson, H.F. Weinberger: Multidimensional nonlinear diffusion arising in population genetics, Adv. Math. 30, (1978), 33-76.
- [3] Z.S. Feng: On explicit exact solutions to the compound Burgers-KdV equation, Phys. Lett. A 293, (2002), 57-66.
- [4] Z.S. Feng: The first-integral method to study the Burgers-Korteweg-de Vries equation, J. Phys. A: Math. Gen. 35, (2002), 343-349.
- [5] R. FitzHugh: Impulses and physiological states in theoretical models of nerve membrane, Biophys. J. (1961) (1-2), 445-466.

- [6] J-H. He, M.A. Abdou: New periodic solutions for nonlinear evolution equations using Exp-function method, Chaos Solitons and Fractals, in press.
- [7] H. Li, Y. Guo: New exact solutions to the Fitzhugh—Nagumo equation, Applied Mathematics and Computation, in press.
- [8] J. D. Murray: Mathematical Biology, Springer 1993, p 304.
- [9] J.S. Nagumo, S. Arimoto, Y. Yoshizawa: An active pulse transmission line simulating nerve axon, Proc. Inst. Radio. Eng. 50 (1962), 2061–2070.
- [10] P. Pelce, J. Sun, C. Langeveld: A Simple model for excitation-contraction coupling in the heart, Chaos, Solitons and Fractals, Vol. 5, No. 314, (1995), 383-391.
- [11] A.D. Polyanin, V.F.Zaitsev: Nonlinear Partial Differential Equations, Chapman&Hall/CRC, 2004.

# ТОЧНИ РЕШЕНИЈА ЗА СЕМЕЈСТВОТО НА FITZHUGH-NAGUMO-ВИДОТ РАВЕНКИ

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### Резиме

Во овој текст се проучува семејството на делумни диференцијални равенки од видот Fitzhugh-Nagumo користејќи го резултатот претставен во [7]. По воведот во проблемот, ги наоѓаме точните решенија на равенката која се применува. Се применува првиот интегрален метод.

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